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Mn4 Single-Molecule-Magnet-Based Polymers of a One-Dimensional Helical Chain and a Three-Dimensional Network: Syntheses, Crystal Structures, and Magnetic Properties

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S Supporting Information

[ABSTRACT:](#page-8-0) Two Mn_4 single-molecule-magnet (SMM)based coordination polymers, $\{[\text{Mn}_4\text{O}(\bar{s}alox)_3(\text{N}_3)_3-\}$ $(DMF)_{2}(H_{2}O)(dpp)]$ ·0.5MeOH}_n (1·0.5MeOH; H₂salox = salicylaldoxime; dpp = 1,3-di-4-pyridylpropane; $DMF = N$, N-dimethylformamide) and $\{[Mn_4O(Me\text{-}salox)_3(N_3)_3]$ $(dpp)_{1.5}$]·1.5Et₂O}_n (2·1.5Et₂O; Me-H₂salox = hydroxyphenylethanone oxime), are self-assembled from $Mn(CIO₄)₂·6H₂O/$ H_2 salox and $Mn(CIO_4)_2·6H_2O/Me-H_2$ salox systems with dpp and NaN_3 in DMF/MeOH , respectively. Both compounds comprise a mixed-valence tetranuclear manganese core, $[Mn^H Mn^H{}_{3}O]^{9+}$, which serves as a building unit for subsequent assembly via oximate and azido ligands. The

flexible dpp ligand links with a Mn₄ unit, leading to the formation of a one-dimensional helical structure in 1·0.5MeOH and a three-dimensional pcu network in 2·1.5Et₂O. The magnetic data analysis shows that antiferromagnetic interactions within the Mn_4 units resulted in S = $\frac{3}{2}$ and $\frac{7}{2}$ ground states for 1·0.5MeOH and 2·1.5Et₂O, respectively. Both compounds show SMM behavior, as evidenced by frequency-dependent out-of-phase signals in alternating-current magnetic susceptibility and magnetic hysteresis loop studies with an energy barrier of $U_{\text{eff}} = 37$ K for $2.1.5Et_2O$.

ENTRODUCTION

Polynuclear manganese clusters have received considerable interest in recent years. Considerable effort has been devoted to understanding the structures and properties of the synthetic models for mimicking the active site of metalloenzymes that contain oxide-bridged manganese cores, such as the wateroxidizing complex of photosystem $II¹$ and catalase enzymes.² Manganese clusters often exhibit large, and sometimes abnormally large, spin values in the gr[ou](#page-8-0)nd state, and combine[d](#page-8-0) with a large anisotropy, some of these species are single-molecule magnets (SMMs) or single-chain magnets (SCMs). SMMs show a significant barrier to magnetization relaxation arising from the combination of a large ground-state spin, S, with an easy-axis anisotropy (negative zero-field-splitting parameter, D).^{3−6} SCMs exhibit a large uniaxial anisotropy and strong intrachain exchange interactions without spin compensation between the [hig](#page-8-0)h-spin magnetic units, with good isolation of the chains that prevents two- $(2D)$ and three-dimensional $(3D)$ ordering.^{7−9} An important future development for SMMs and SCMs is the discovery of synthetic schemes that can yield new co[mpou](#page-8-0)nds and families of related compounds with large spins and/or significant magnetoanisotropies.¹⁰ It is known that the Mn ^{III}ion Jahn−Teller (JT) elongation plays an importance role in SMM and SCM properties a[nd](#page-8-0) that the orientations of JT elongation significantly affect the energy barrier.¹¹ Thus, many current routes for producing SMMs and SCMs involve Mn^{III}containing compounds to exploit the associated [lar](#page-8-0)ge single-ion anisotropy.^{9a,b,10}

More recently, the use of manganese clusters as building blocks for [the c](#page-8-0)onstruction of coordination polymers as onedimensional $(1D)$,¹² 2D,¹³ and, in very few examples, $3D$ ¹⁴ structures with interesting magnetic properties was also investigated. Theo[ret](#page-8-0)icall[y,](#page-8-0) such an arrangement of molecu[lar](#page-8-0) magnets in a crystalline material would enable magnetization probing without the interference of magnetic ordering, allowing researchers to study the exact nature of the retention of magnetization for S MMs.^{14a,15} A synthetic methodology that has proven to be successful in the synthesis of such compounds is the reaction of linking li[gands](#page-8-0) with chelating coligands, which

Received: July 3, 2012 Published: December 6, 2012 may subserve the formation of manganese clusters. Salicylaldoxime-containing manganese compounds are considered to be good candidates as building units for these materials because they usually comprise a $[Mn_3O]^{7+}$ triangle core, which is magnetically interesting because of the pronounced magnetic anisotropy from its almost near-parallel JT distortion axes in high-spin Mn^{III} ions. Some of these compounds have been reported as $SMMs^{16}$ and as building blocks of a 1D chain.¹⁷

Using salicylaldoxime and its derivatives (H₂salox and Me- $H₂$ sal[o](#page-9-0)x) with azido ligands and 1,3-di-4-pyridylpropane (d[pp](#page-9-0)) is attractive for the above applications. The combination of salicylaldoxime and azido ligands is known in cluster coordination chemistry especially for $SMMs$,¹⁸ but it has few uses in coordination polymers.^{17a,19} The present study used these ligands to prepare two new compounds, $\{[Mn_4O(salox)_3(N_3)(DMF)_2\}$ $(H_2O)(dpp)]$ ·0.5MeO[H}](#page-9-0)_{[n](#page-9-0)} (1·0.5MeOH; H₂salox = salicylaldoxime; dpp = 1,3-di-4-pyridylpropane; DMF = N,N-dimethylformamide) and $\{[Mn_4O(Me\text{-}salox)_3(N_3)_3(dpp)_{1.5}]\cdot 1.5Et_2O\}_n$ $(2.1.5Et₂O; Me-H₂salox = hydroxyphenylethanone oxime).$ The former adopts an 1D helical structure, and the latter shows a 3D pcu network. Both compounds are constructed from an unprecedented Mn_4 topology building unit. These two compounds not only show interesting topologies but also exhibit SMM behaviors resulting from the significant magnetic anisotropy of the $Mn₄$ unit. The detailed syntheses, structures, and magnetic properties of these compounds are described below.

EXPERIMENTAL SECTION

Synthesis. All solvents and reagents where used as-received; no purification was necessary. All reactions were performed under aerobic conditions. Caution! Perchlorate compounds are potentially explosive; such compounds should be synthesized, used in small quantities, and treated with the utmost care at all times. The derived oximes were prepared via a condensation reaction, as previously reported. 20

 ${[(Mn_4O(salox)_{3}(N_3)_{3}(DMF)_{2}(H_2O)(dpp)]}\cdot 0.5MeOH\}$ _n (1.0.5MeOH). $Mn(CIO₄)$ ₂·6H₂O (181 mg, 0.50 mmol) a[nd](#page-9-0) NaN₃ (64.5 mg, 0.99 mmol) were dissolved in a 1:1 MeOH/DMF mixture (20 mL). To this was added H₂salox (68.3 mg, 0.50 mmol) and NEt₃ (50.4 mg, 0.50 mmol), resulting in a deep-green solution. After stirring for 5 min, dpp (100.1 mg, 0.50 mmol) was added, followed by stirring for 30 min, during which time the solution turned dark green. The vapor diffusion of Et_2O into the solution slowly produced crystals. After 1 week, well-formed deep-green crystals appeared. The sample for crystallographic study, identified as 1·0.5MeOH, was maintained in contact with the mother liquor to avoid solvent loss. The yield was 60% (based on Mn). Anal. Calcd (found) for $C_{40.5}H_{47}Mn_4N_{16}O_{10.5}$ (1·0.5MeOH): C, 47.46 (47.18); H, 4.13 (4.12); N, 19.56 (19.63). IR data (KBr disk, cm[−]¹): 3422 (br, s), 2928 (m), 2075 (vs), 2059 (vs), 1655 (s), 1596 (s), 1541 (m), 1439 (s), 1386 (w), 1327 (m), 1284 (s), 1202 (m), 1152 (w), 1122 (w), 1107 (w), 1042 (s), 1068 (s), 918 (s), 819 (w), 760 (s), 679 (vs), 647 (s), 465 (m).

 ${[(Mn_4O(Me-salox)_{3}(N_3)_{3}(dpp)_{1.5}]\cdot 1.5Et_2O]}$ (2.1.5Et₂O). The procedure was the same as that employed for compound 1, except that an excess of Me-H₂salox (75.5 mg, 0.50 mmol) was added instead of H2salox. Deep-green crystals were obtained. The yield was 68% (based on Mn). The sample used for crystallography, identified as $2.1.5Et₂O$, was maintained in contact with the mother liquor to avoid solvent loss. Dried solid was hygroscopic and analyzed as 2.2H₂O. Anal. Calcd (found) for $C_{43.5}H_{36}Mn_4N_{16}O_{11}$ (2·2H₂O): C, 45.72 (46.27); H, 4.06 (4.68); N, 18.39 (18.07). Selected IR data (KBr disk, cm[−]¹): 3433 (br, s), 2930 (m), 2057 (vs), 1653 (s), 1612 (s), 1595 (s), 1569 (m), 1527 (m), 1473 (w), 1434 (s), 1384 (w), 1315 (s), 1281 (w), 1247 (m), 1134 (w), 1052 (m), 1021 (m), 971 (s), 864 (m), 806 (w), 759 (m), 678 (vs), 650 (m), 607 (m), 571 (m), 545 (w), 519 (w), 477 (w), 438 (m).

X-ray Crystallography. The single crystals of 1·0.5MeOH and $2.1.5Et₂O$ were mounted on the tip of a glass fiber with dimensions of 0.50 \times 0.35 \times 0.03 and 0.22 \times 0.20 \times 0.20 mm³, respectively.

Intensity data were collected at $150(2)$ K within the limits of 1.49° < θ < 27.50° for 1.0.5MeOH and 1.44° < θ < 27.50° for 2.1.5Et₂O using a Nonius Kappa CCD diffractometer equipped with graphitemonochromatized Mo K α radiation ($\lambda = 0.71073$ Å). The structures were solved by direct methods and refined by the full-matrix leastsquares method on F^2 values using the SHELX-97²¹ program package. All non-hydrogen atoms were refined anisotropically, whereas the hydrogen atoms were placed in ideal, calcula[ted](#page-9-0) positions, with isotropic thermal parameters riding on their respective carbon atoms. In 1·0.5MeOH, one of the coordinated DMF molecules is disordered over two positions with site occupation factors (SOFs) of 0.55 and 0.45, and the solvated MeOH is disordered with a SOF of 0.5. In $2·1.5Et₂O$, $Et₂O$ is disordered over two positions with a SOF of 0.5. The experimental details for X-ray data collection and refinements of 1.0.5MeOH and $2.1.5Et₂O$ are presented in Table 1.

Physical Measurements. IR spectra were recorded for the solid state (KBr pellets) on a Nicolet Magna 550 Fourier transform IR spectrometer in the range of 4000–400 cm⁻¹. Elemental analyses were carried out using an Elementar Vario EL III analyzer. Thermogravimetric analyses (TGA) were performed under nitrogen with a Shimadzu TGA-50 analyzer. Powder X-ray diffraction (PXRD) data were recorded on a Shimadzu XRD-7000S X-ray diffractometer at 30 kV and 30 mA with Cu K α radiation ($\lambda = 1.5406$ Å), with a step size of 0.028 in θ and a scan speed of 1 s step⁻¹. Variable-temperature direct-current (dc) and alternating-current (ac) magnetic susceptibility measurements were collected on microcrystalline samples, restrained in eicosane to prevent torquing, on a Quantum Design MPMS-XL SQUID and a PPMS magnetometer equipped with 7.0 and 9.0 T magnets, respectively, operated in the range of 1.8−300.0 K. Diamagnetic corrections were estimated from Pascal's constants 22 and subtracted from the experimental susceptibility data to obtain the molar paramagnetic susceptibilities of the compounds.

■ RESULTS [A](#page-9-0)ND DISCUSSION

Synthesis. The salicylaldoxime group (H_2salox) has been used as a chelating ligand in metal cluster chemistry.²³ It has been observed in its singly deprotonated (Hsalox[−]) form as a

bidentate chelate with the protonated N–OH not bound,²⁴ but it more commonly occurs in its doubly deprotonated (salox^{2−}) form as a tridentate chelate.²⁵ The variety of binding modes of H_2 salox may allow the synthesis of new nuclearities and topologies. The treatment of a MeOH[/D](#page-9-0)MF mixture solution of $Mn(CIO₄)₂·6H₂O$, NEt₃, NaN₃, H₂salox, and dpp in a 1:1:1:1:1 ratio led to a darkgreen reaction mixture. The residue was slowly diffused with $Et₂O$ to yield dark-green-brown triangular crystals of 1·0.5MeOH in 60% yield. When the reaction process of compound 1 was carried out with equivalent Me-H₂salox instead of H₂salox, compound $2.1.5Et$, O was obtained in a yield of 68%. The reactions of compounds 1 and 2 are summarized in eqs 1 and 2, respectively.

$$
16Mn(ClQ4)2·6H2O + 12H2salox + 12NaN3 + 8DMF+ 4dpp + 2MeOH + 3O2→ 41·0.5MeOH + 32ClQ4- + 94H2O+ 20H+ + 12Na+
$$
 (1)

 $16Mn(ClO_4)_2.6H_2O + 12Me-H_2salox + 12NaN_3 + 6dpp$

+
$$
6Et_2O + 3O_2
$$

\n $\rightarrow 4\{[Mn_4O(Me-salox)_3(N_3)_3(dpp)_{1.5}] \cdot 1.5Et_2O\}_n$
\n+ $32ClO_4^- + 94H_2O + 20H^+ + 12Na^+$ (2)

These reactions involve manganese oxidation, undoubtedly by O_2 under the prevailing basic conditions, and have been balanced accordingly. NEt_3 ensures basic conditions and acts as a proton acceptor; in its absence, longer reaction times (>24 h) are required to get a significant dark-red coloration, and the yields of isolated 1.0.5MeOH and $2.1.5Et₂O$ are much lower (<5%). On the other hand, >2 equiv of NEt₃ gave insoluble amorphous precipitates that were probably manganese oxides or oxo/hy droxides. An increase in the amount of H_2 salox or Me-H2salox to 3 equiv (or more) led to the isolation of 1.0.5MeOH and $2.1.5Et₂O$ in low yields of <10%.

Description of the Crystal Structures. Single crystals of compounds $1.0.5$ MeOH and $2.1.5$ Et₂O suitable for singlecrystal X-ray crystallography were obtained. Labeled plots of crystal structures and packing diagrams for both compounds are shown in Figures 1−5. The selected bond distances and angles are listed in Table 2.

1·0.5MeOH. X-[ra](#page-3-0)[y s](#page-4-0)tructural analysis shows that 1·0.5MeOH crystallized in monoclinic space group $C2/c$ and that the asymmetric unit contains the whole Mn_4 unit and half of MeOH. As shown in Figure 1, compound 1·0.5MeOH has a covalently bonded 1D structure, which was built from Mn_4 clusters connected by one d[pp](#page-3-0) ligand. The 1·0.5MeOH molecule does not have any crystallographic symmetry. Bond-valence-sum $(BVS)^{26}$ calculations (Table 3) and close examination of the structural parameters agree with a valence-trapped description of

Figure 1. ORTEP drawing of the asymmetric unit of 1·0.5MeOH with thermal ellipsoids set at 30% probability. Hydrogen atoms and solvated molecules have been omitted for clarity (right). The 1D chain structure of 1·0.5MeOH is viewed along the a axis.

Mn1−Mn3 being Mn^{III} and Mn4 being Mn^{II}. The Mn^{III} centers of Mn1 and Mn2 are six-coordinated, each of them clearly possessing a JT distortion in the form of axis elongation along N1−Mn1−N14 and N7−Mn2−O10, whereas Mn3 adopts a five-coordinated geometry, which is regarded as a square pyramid $(\tau < 0.001)$, in which the nitrogen atom of an end-on azide is located in the apical position.²⁷ As expected, the \overline{I} axes are located so as to avoid Mn−O2[−] bonds, which are the shortest

Figure 2. Structure of a double-helical chain assembly of 1. Figure 3. ORTEP drawing of a Mn₄ cluster in 2·1.5Et₂O. Hydrogen atoms and solvated molecules have been omitted for clarity.

and strongest in the molecule $(\langle 1.9 \text{ Å} \rangle)$, and the JT axes at the two Mn^{III} ions in the compound are nearly parallel. The structure of the Mn₄ cluster can be described as a distorted triangular pyramid, whose basal plane is constructed by a μ -O^{2−}, which lies ∼0.118 Å below the plane of the three manganese ions and whose centered $[\mathrm{Mn}^{\mathrm{III}}_{3}\mathrm{O}]$ triangle and apical is occupied by a $\mathrm{Mn}^{\mathrm{II}}$ ion. Each Mn^{III} ion of the triangle is bridged to the apical Mn^{II} by three end-on azido ligands. Each edge of the triangle is bridged by

Figure 4. Structure drawing of the $(Mn_4)_2$ unit in 2.

Figure 5. Schematic representation of a 3D pcu framework in 2, including chemical structures of the components.

Table 3. BVS Calculation Results of Manganese Centers for Compounds 1.0 .5MeOH and 2.1 .5Et₂O

		Mn1	Mn2	Mn3	Mn4
Compound 1.0.5MeOH					
	Mn^{II}	3.24	3.33	3.31	1.84
	Mn ^{III}	2.99	3.07	3.05	1.70
	Mn^{IV}	2.93	3.01	2.99	1.67
Compound 2.1.5Et,O					
	Mn^{II}	1.65	3.19		
	Mn^{III}	1.53	2.94		
	$Mn^{\rm IV}$	1.5	2.89		

a dianionic oximate group of a salox $^{2-}$ ligand in the $\eta^1:\eta^1:\mu^1:\mu_2$ mode, whose each deprotonated hydroxyl group is bound terminally to a Mn^{III} ion to form a Mn−O−N−Mn chelating ring. The chelating rings are slightly distorted from planarity with torsion angles of 14.81°, 0.57°, and 15.60° for Mn1−O5−N2− Mn2, Mn2−O7−N3−Mn3, and Mn3−O3−N1−Mn1, respectively. The peripheral ligations in the $Mn₄$ unit are provided by one water molecule, two DMF molecules, and one dpp ligand. The flexible dpp ligand adapts a trans−gauche conformation and connects to Mn1 $'$ in the neighboring Mn₄ unit, forming a 1D helical structure (Figure 2) along the b axis. Moreover, the adjacent left- and right-handed helical chains are cross-linked into a double-helical chain struc[tu](#page-3-0)re by the hydrogen bonding of solvated MeOH.

 $2.1.5Et₂O$. Single-crystal X-ray diffraction analysis reveals that compound $2.1.5Et₂O$ crystallizes in rhombohedral space group $R\overline{3}c$. Similar to compound 1, the compound contains a mixed-valence $Mn^{II}Mn^{III}$ ₃ cluster (Figure 3), the structural geometry of each Mn_4 cluster in which the oxidation state of each metal center is evidenced by BVS ([Ta](#page-3-0)ble 3) and JT elongation axes. The Mn_4 cluster adopts a distorted trigonalpyramidal geometry with a basal plane of a $[Mn_{3}^{III}O]$ triangle and an apex occupied a Mn^{II} ion. The Mn^{II} ion of $Mn1$ is sixcoordinated with a distorted octahedral geometry, and the Mn^{III} center of Mn2 adopts a five-coordinated geometry, which is regarded as a distorted square pyramid ($\tau = 0.31$), in which the nitrogen atom of the end-on azide is located in the apical position. The Me-salox²⁻ ligand bridges the edge of the [Mn^{III}₃O] triangle with a resulting Mn2−O3−N4−Mn2′ torsion angle of 43.13°. In compound $2.1.5Et₂O$, two $Mn₄$ clusters are connected in the face-to-face arrangement of their $[Mn_{3}^{11}]$ to form a (Mn_4) subunit through the weak interaction of Mn2…O3′ with a distance of 2.736 Å (Figure 4). The $(Mn_4)_2$ subunits are further linked by three dpp ligands (Mn1−Mn1′) in a trans−trans conformation, forming a 3D framework. If the $(Mn_4)_2$ units linked by three dpp groups are viewed as one node, compound $2.1.5Et₂O$ can be considered as a six-connected pcu topology structure with a 2-fold interpenetrating 3D architecture with another independent equivalent framework (Figure 5).

TGA and PXRD. The thermal stabilities of 1·0.5MeOH and $2.1.5Et₂O$ were examined. The results of TGA are in good agreement with the crystallographic observations. A TGA trace of 1 shows that water removal continues in several steps, resulting in a total weight loss of 15.4% in the temperature range of 30−200 °C, corresponding to the loss of a guest MeOH molecule, two coordinated DMF molecules, and one water molecule (calcd, 15.8%). Compound 2 is stable up to 150 °C. Decomposition of 1 and 2 started at 200 and 150 °C, respectively (Figure S1 in the Supporting Information).

PXRD patterns were obtained for compounds 1·0.5MeOH and $2.1.5Et₂O$. The PXRD pat[terns of the two compou](#page-8-0)nds are in good agreement with the patterns simulated from the respective single-crystal data, confirming the bulk purity of the two samples (Figures S2 and S3 in the Supporting Information).

Magnetic Properties. Solid-state, variable-temperature magnetic susceptibility measurement[s were performed on mic](#page-8-0)rocrystalline samples of compounds $1.0.5$ MeOH and $2.2H₂O$, all suspended in eicosane to prevent torquing. The dc magnetic susceptibility data of compounds $1.0.5$ MeOH and $2.2H₂O$ were collected in the 2.0−300 K range in a 1.0 kG magnetic field and are illustrated in Figures 6a and 7a, respectively. The value of $\chi_{\rm M}$ T for 1·0.5MeOH steadily decreases from 11.4 emu K mol⁻¹ at 300 [K](#page-5-0) to 5.8 emu K mol⁻¹ at 14 K, below which the $\chi_{\rm M}T$ value decreases rapidly to 3.4 [em](#page-5-0)u mol[−]¹ K at 1.8 K. The values at 300 K are slightly less than the calculated spin-only value of 13.4 emu K mol⁻¹ for a Mn^{II}Mn^{III}₃ compound with noninteracting metal centers with $g = 2.0$, suggesting antiferromagnetic interactions. The monotonically decreasing $\chi_{\text{M}}T$ with temperature and the resultant low value of $\chi_{\rm M}T$ at 2.0 K are indicative of a small ground state for compound 1·0.5MeOH. In order to characterize the coupling interactions between manganese centers within the Mn_4 unit of 1.0.5MeOH, the magnetic susceptibility data were fit to the appropriate theoretical expression. Considering the topologies and symmetry of the $Mn₄$ cluster, a three-J model was introduced (Figure 6b). In the model, couplings are characterized as Mn^{III}-Mn^{III} through μ_3 -O^{2−} and oximate bridgings ([M](#page-5-0)n1–Mn2 and Mn1–Mn3 for J₁

Figure 6. (a) Plot of $\chi_{\rm M}T$ versus T for microcrystalline samples of compound 1·0.5MeOH. The solid lines represent a least-squares fit of the data in the region 10−300 K to the van Vleck equation. (b) Diagram showing the definitions of the atom number and magnetic exchange parameters for compound 1·0.5MeOH.

and Mn2−Mn3 for J_2 based on the torsion angle) and Mn^{II}− Mn^{III} through $\mu_{1,1}N_3^-$ bridging (Mn1−Mn4, Mn2−Mn4, and Mn3−Mn4 for J_3). The spin Hamiltonian that describes the isotropic magnetic exchange interaction is given as

$$
H = -2J_2(S_2S_3) - 2J_1(S_1S_2 + S_1S_3)
$$

- 2J_3(S_1S_4 + S_2S_4 + S_3S_4) (3)

For a compound with three $S = 2$ and one $S = \frac{5}{2}$ interacting ions, the overall degeneracy of this spin system is 95, which make up a total of 9 possible states with S_{T_2} the total spin of a $Mn^H Mn^H$ ₃ cluster, ranging from 0 to $17/2$. A theoretical expression for the molar paramagnetic susceptibility (χ_M) versus temperature (T) was derived using the van Vleck equation and assuming an isotropic g value. This expression was used to fit the experimental χ_M versus T data for compound 1.0.5MeOH; the parameters varied were J and g . The data below 10 K was omitted in the fitting because zero-field-splitting and Zeeman effects are likely to have an effect on the susceptibilities in this temperature range. The results of fitting the experimental data are shown as solid lines in Figure 6a, with final parameters of $g = 1.96$, $J_1 = -3.5$ cm⁻¹, $J_2 = -9.3$ cm⁻¹, and $J_3 = -0.2$ cm⁻¹ . The temperature-independent paramagnetism (TIP) was held constant at 800 \times 10⁻⁶ emu mol⁻¹. This set of parameters results in a ground state of $S = \frac{3}{2}$ for 1.0.5MeOH, with a very closelying first excited state $(S = \frac{1}{2})$ located only 0.8 cm⁻¹ above and the second excited state $(S = \frac{5}{2})$ at 1.2 cm⁻¹ above.

For compound $2.2H_2O$, $\chi_M T$ slightly decreases from 12.2 emu K mol⁻¹ at 300 K to reach a value of 11.7 emu K mol⁻¹ at 50 K. Upon cooling, $\chi_M T$ decreases again to 8.6 emu K mol⁻¹ at

Figure 7. Plot of $\chi_M T$ versus T for microcrystalline samples of compound $2.2H₂O$. The solid lines represent a least-squares fit of the data in the region 10−300 K to the van Vleck equation. (b) Diagram showing the definitions of the atom number and magnetic exchange parameters for compound $2.2H₂O$.

2.0 K. The value of $\chi_{\rm M}T$ at 300 K is slightly lower than the calculated spin-only value of 13.4 emu K mol⁻¹ for a Mn^{II}Mn^{III}₃ compound of noninteracting metal centers with $g = 2.0$, suggesting that weak antiferromagnetic couplings dominate the overall intraunit exchange interactions within the compound. The decreasing $\chi_M T$ at low temperature is likely due to the zero-field-splitting, Zeeman effects and/or intermolecular interactions of the high-spin ground state. Because of the topological complexity of the $(Mn₄)₂$ unit, it is not possible to determine the individual pairwise manganese exchange interactions using the Kambe method.²⁸ Thus, analysis of the magnetic interactions in $2.2H₂O$ was simplified into a $Mn₄$ cluster, which ignores the magneti[c i](#page-9-0)nteraction between two staggered $Mn₄$ clusters through a weak interaction of Mn2···O3′ at a distance of 2.736 Å. Considering the C_{3v} symmetry of the Mn₄ cluster in $2.2H₂O$, the coupling scheme is shown in Figure 7b. On the basis of the nature of the Mn−Mn distances, a two-J model was used, where J_1 and J_2 characterize the coupling within $Mn^{III}-Mn^{III}$ through μ_3 -O²⁻ and oximate bridging and within $\text{Mn}^{\text{II}}\text{--}\text{Mn}^{\text{III}}$ through $\mu_{1,1}$ - N_3^{-} bridging, respectively. The spin Hamiltonian is given as

$$
H = -2J_1(S_1S_2 + S_2S_3 + S_1S_3) - 2J_2(S_1S_4 + S_2S_4 + S_3S_4)
$$
\n(4)

The data below 10 K were omitted in the fitting because zerofield-splitting, Zeeman effects and interunit interactions are likely to have an effect on the susceptibilities in this temperature range. The results of fitting the experimental data are shown as solid lines in Figure 7a, with final parameters of $g = 1.91$, $J_1 = 0.7$ cm⁻¹, and $J_2 = -0.6$ cm⁻¹. The TIP was held constant at 800 \times 10⁻⁶ emu mol⁻¹. This set of parameters

results in a ground state of $S = \frac{7}{2}$ for 2·2H₂O, with a closelying first excited state $(S = \frac{9}{2})$ located only 5.4 cm⁻¹ above and the second excited state $(S = \frac{5}{2})$ at 11.3 cm⁻¹ above.

The values obtained for the exchange parameters of the two compounds are comparable to those for a similar pathway in other reported compounds. In compound 1, the intratriangular antiferromagnetic interactions $(J_1 \text{ and } J_2)$ may result from bridges of μ_3 -O^{2−} and the small torsion angles of salicylaldoximate (below 31°). These are consistent with reported compounds that contain a salicylaldoxime-based $[Mn_3O]^{7+}$ triangle.16c−f,29 The intratriangular interactions in compound 2 (I_1) are ferromagnetic and result from the twisting of Mn−N−[O](#page-9-0)[−](#page-9-0)[M](#page-9-0)n arrangements with an angle of 43.11°, which is due to the small overlapping of the d_z^2 magnetic orbital for manganese ions³⁰ and consistent with reported results.^{16a,d,29b,31} The small J_1 value in compound 2 may be attributed to the antiferromagne[tic](#page-9-0) interaction between two staggered $Mn₄$ clusters. The interactions between neighboring Mn^{II} and Mn^{III} ions bridged by end-on N_3^- are antiferromagnetic $(J_3$ in 1 and J_2 in 2), which produces the large bridging angles of the azide ions (121.3−127.3° for 1 and 121.7° for 2).32

The S values of the $\rm Mn_4$ clusters in 1·0.5MeOH and $\rm 2\cdot 2H_2O$ were confirmed from the variable-te[mp](#page-9-0)erature and variablefield magnetization (M) measurements in the 2.0−4.0 K and 1−7 T ranges, respectively (Figure S4 and S5 in the Supporting Information). The data were fitted using the program $ANISOFIT, ³³$ assuming that only the ground state i[s populated.](#page-8-0) [However, an](#page-8-0) acceptable fit could not be obtained using the data collec[ted](#page-9-0) over the whole field range, which is a common problem because of low-lying excited states if some show an S value greater than that of the ground state. A common solution is to only use data collected with low fields (below 1.0 T).³⁴ However, it was still not possible to obtain a satisfactory fit assuming that only the ground state is populated in t[his](#page-9-0) temperature range. This suggests that both compounds possess particularly low-lying excited states that are populated even at these relatively low temperatures. This is consistent with fitting results for the dc magnetic measurements.

The measurements of the ac magnetic susceptibility were performed on microcrystalline samples for compounds 1·0.5MeOH and 2·2H2O in the temperature range of 1.8−5 K in a zero dc field and a 3.5 G ac field oscillating at frequencies in the range of 100−10000 Hz. The ac susceptibility data for 1·0.5MeOH are presented in Figure 8, which shows the inphase ac susceptibility χ_{M} ['] (as χ_{M} [']T) versus T (Figure 8, top) and the out-of-phase ac susceptibility χ_M'' versus T (Figure 8, bottom) plots. The sloping $\chi_M' T$ versus T plot is strongly indicative of a population of low-lying excited states because occupation of only the ground state would give an essentially temperature-independent value. At lower temperatures (below 4 K), the in-phase signal decreases and a frequency-dependent χ_M ["] signal appears, suggesting slow relaxation of the SMM behavior. However, the peak maxima clearly lie at temperatures below 1.8 K, the operating limit of our instrument. Extrapolation of $\chi_{\rm M}^{\phantom i}$ data down to 0 K gives ~1.7 emu K mol $^{-1}$, indicating an $S = \frac{3}{2}$ ground state with $g \sim 1.95$, which is in good agreement with the dc magnetization fits.

The ac magnetic susceptibility data of compound $2.2H₂O$ are shown in Figure 9. The χ_{M} 'T values of 2.2H₂O are essentially temperature-independent down to ∼4 K, below which they decrease because of slow relaxation. Extrapolate χ_{M} ^T to a value in the of 9.5 emu K mol⁻¹ range at 0 K, suggesting a ground state of $S = \frac{7}{2} \pm 1$, which is in good agreement with dc

Figure 8. Plots of (top) χ_M/T and (bottom) χ_M'' versus T for a microcrystalline sample of compound 1·0.5MeOH in a 3.5 G ac field. The data were collected in an ac field oscillating at the indicated frequency.

Figure 9. Plots of (top) $\chi_M' T$ and (bottom) χ_M'' versus T for a microcrystalline sample of compound $2.2H₂O$ in a 3.5 G ac field. The data were collected in an ac field oscillating at the indicated frequency.

studies.^{32b} The out-of-phase χ_{M} " signals increased with decreasing temperature; they reached a maximum value at 2.6 K [in 1](#page-9-0)0000 Hz and then approached zero, which was frequency dependent. This frequency dependence of the ac signals, which indicates superparamagnetic behavior in compound 2, is caused by the inability of the compound to relax quickly enough to keep up with the oscillating field at these temperatures. The peak temperatures of χ_M " can be measured by the parameter $\phi = (\Delta T_p/T_p)/\Delta(\log f) = 0.13$, which is close to the range of normal superparamagnets.^{35,36}

In order to confirm the SMM properties, magnetization versus applied dc field data down to 0.04 K were collect[ed o](#page-9-0)n single crystals of $1.0.5$ MeOH and $2.1.5$ Et₂O using a micro-SQUID apparatus.³⁷ Magnetization versus dc field scans for single crystals of compounds $1.0.5$ MeOH and $2.1.5$ Et₂O at

Figure 10. Magnetization hysteresis loops for a single crystal of 1·0.5MeOH: (top) temperature dependence at a fixed scan rate of 0.14 T/s; (bottom) scan-rate dependence at a fixed temperature of 0.03 K. M is normalized to its saturation value, M_s .

Figure 11. Magnetization hysteresis loops for a single crystal of $2.1.5Et₂O$: (top) temperature dependence at a fixed scan rate of 0.14 T/s ; (bottom) scan-rate dependence at a fixed temperature of 0.03 K. M is normalized to its saturation value, M_s .

various temperatures and scan rates are shown in Figures 10−12, respectively. For compound 1·0.5MeOH, the corresponding magnetization responses are at 0.04−1.3 K and a fixed field sweep rate of 0.14 T/s and magnetization responses at

Figure 12. (a) Magnetization decay versus time plots for a single crystal of $2.1.5Et₂O$ at the indicated temperatures and (b) relaxation time (τ) versus $1/T$ plot for 2·1.5Et₂O constructed using dc magnetization decay data. The dashed line is the fit of the thermally activated region to the Arrhenius equation. See the text for the fit parameters.

0.28−0.002 T/s field sweep rates and a constant temperature of 0.03 K. Below 0.4 K, hysteresis loops are observed, whose coercivities increase with increasing field sweep rate and decreasing temperature, indicative of SMM behavior. For compound $2.1.5Et₂O$, the corresponding magnetization responses are at $0.03-1.8$ K and a fixed field sweep rate of 0.14 T/s, which shows stepwise magnetization and large coercivities, indicating the occurrence of a quantun tunneling mechanism and significant magnetic anisotropy. Taking into account the crystal structure, we suggest that this large magnetic anisotropy might come from the elongated axes (N1−Mn2…O3) of the six Mn^{III} ions of the two-faced $Mn₄$ units. Again, the coercivities of the hysteresis loops increase with increasing field sweep rate and decreasing temperature, as expected for the superparamagnet-like behavior of SMMs. Analyzing these hysteresis loops at several field scan rates in more detail, three steps at $H = 0$, 0.3, and 0.9 T are observed. The steps at $H = 0$ and 0.9 T are rather sharp, whereas the step at 0.3 T is broad. These steps are probably due to resonant quantum tunneling between ground- and/or excited-state energy levels of the coupled $[Mn_4]_2$ dimer units. In order to determine the effective barrier (U_{eff}) to magnetization relaxation, relaxation rate versus T data at various temperatures down to 0.04 K were obtained from a magnetization decay versus time study (Figure 12). The resulting relaxation rates $(1/\tau)$ versus T from ac and dc decay data were used to construct an Arrhenius plot, shown as τ versus $1/T$ in Figure 12, based on the Arrhenius relationship. The fit to the thermally activated region gave $\tau_0 = 4.3 \times 10^{-10}$ s and $U_{\text{eff}} = 37$ K. This effective energy barrier and τ_0 are close to those given for previously reported SMMs.3[−]6,11,12,16 Below ∼0.1 K, the relaxation becomes independent of the temperature, consistent with relaxation only via a quantun tunn[eling me](#page-8-0)[ch](#page-9-0)anism through the anisotropy barrier.

■ CONCLUSION

Two new coordination polymers based on mixed-valence tetranuclear manganese Mn^{III} Mn^{II} SMMs, 1.0.5MeOH and 2.1.5Et₂O, were self-assembled using the flexible 1,3-di(4-pyridyl)propane (dpp) ligand with coligands of salicylaldoxime and azide. The flexible moiety of dpp modulated 1·0.5MeOH and $2.1.5Et₂O$, forming a 1D helical chain structure and a 3D framework, respectively. The dc magnetic susceptibility measurements indicate that compounds $1.0.5$ MeOH and $2.1.5$ Et₂O exhibit intramolecular antiferromagnetic exchange and possess $S = \frac{3}{2}$ and $\frac{7}{2}$ ground states, respectively, together with significant magnetic anisotropy, resulting in a small barrier to magnetization relaxation, indicating SMM behavior in both compounds. The results demonstrate that the ligand flexibility of the 1,3-di(4-pyridyl)propane ligands with coligands of salicylaldoxime and azide is a useful motif in the design and synthesis of new multidimensional cluster-based magnetic materials.

■ ASSOCIATED CONTENT

S Supporting Information

TGA results (Figure S1), PXRD data (Figures S2 and S3), reduced magnetization data (Figures S4 and S5), and crystallographic files in CIF format. This material is available free of charge via the Internet at http://pubs.acs.org.

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The authors declare no competing financial interest.

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